

Solutions Exam MOL040, Thermodynamics 2

23 January 2019

All 16 subproblems have equal weight for the final grade.
Sometimes a **NOTE** is added, just meant as an explanation or advise.

Problem 1

- a) Note, that the answers below to part a) are not unique, in some cases overcomplete and are meant to be indicative.
- Triple point: A point in a phase diagram where three different phases are in mutual equilibrium.
 - First order phase transition: A transition between two phases for which the first derivative of the chemical potential with respect to the temperature is discontinuous. Consequently the change in volume, enthalpy, entropy will be non-zero.
 - Regular solution: A solution for which the excess enthalpy on mixing $\Delta H^E \neq 0$, but the excess entropy $\Delta S^E = 0$.
 - Molality: The amount of a component i in a mixture in terms of its amount in number of moles divided by the mass of the solvent in kg. $m_i = \frac{n_i}{\text{kg of solvent}}$.
 - pH: $\text{pH} \equiv -\log a(\text{H}^+)$, where $a(\text{H}^+)$ is the activity of the H^+ ions in the solution.

b)

$$dG = VdP - SdT + \mu_A dn_A + \mu_B dn_B + Edq$$

$$\left(\frac{\partial \mu_A}{\partial q}\right)_{P,T,n_A,n_B} = \left(\frac{\partial E}{\partial n_A}\right)_{P,T,n_B,q}$$

- c) Partial molar value of X , with respect to component i at constant temperature and pressure.
 X : any function, e.g. volume, Gibbs free energy, etc.
 i : component i
 n : amount/number of moles
 P : pressure of the system/mixture
 T : temperature of the system/mixture

d) Use the Boltzmann definition of entropy

$$S = k \ln W, \quad \text{so} \quad \Delta_{trs} S = k \ln W' - k \ln W = k \ln \frac{W'}{W}, \quad \text{or} \quad W' = W \exp \left[\frac{\Delta_{trs} S}{k} \right].$$

At the phase transition (Clapeyron equation)

$$\frac{dP}{dT} = \frac{\Delta_{trs} S}{\Delta_{trs} V}, \quad \text{so} \quad \Delta_{trs} S = \left(\frac{dP}{dT}\right) \Delta_{trs} V.$$

Combining this with the earlier result, we find

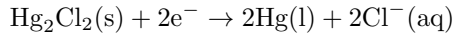
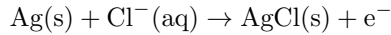
$$W' = W \exp \left[\frac{1}{k} \left(\frac{dP}{dT}\right) \Delta_{trs} V \right].$$

For a transition from solid to gas, the value of $\left(\frac{dP}{dT}\right)$ is negligible for the solid, while for the gas phase it follows from the equation of state $PV = nRT$, so $\left(\frac{dP}{dT}\right) \approx \left(\frac{dP_g}{dT}\right) = \frac{n_g R}{V_g}$, so we get ($nR = Nk$)

$$W' \approx W \exp \left[\frac{n_g R}{k} \frac{\Delta_{trs} V}{V_g} \right] = W \exp \left[N_g \frac{V_g - V_s}{V_g} \right] = W \exp \left[N_g \left(1 - \frac{V_s}{V_g} \right) \right].$$

Problem 2

a)



b)

$$\Delta_r G = -\nu F E = -2 \cdot 9.648 \cdot 10^4 \cdot 0.046 = -8.876 \text{ kJ/mol.}$$

c) At constant pressure ($dP = 0$):

$$dH = dQ + VdP + Edq = dQ + Edq.$$

Reversibly, so

$$\Delta_r H = Q^{rev} + E\Delta q = Q^{rev} + W^{max} = Q^{rev} - \Delta_r G = 11.08 \text{ kJ/mol, so } Q^{rev} = 11.08 - (-8.876) = 19.96 \text{ kJ/mol.}$$

d)

$$dG = VdP - SdT + Edq, \quad \text{so} \quad \left(\frac{\partial E}{\partial T}\right)_{P,q} = -\left(\frac{\partial S}{\partial q}\right)_{P,T} = -\frac{1}{T} \left(\frac{\partial Q^{rev}}{\partial q}\right)_{P,T}.$$

$$\left(\frac{\partial E}{\partial T}\right)_{P,q} = -\frac{1}{T} \frac{dQ^{rev}}{dq} = -\frac{1}{298} \cdot \frac{19.96 \cdot 10^3}{-2 \cdot 9.648 \cdot 10^4} = 3.47 \cdot 10^{-4} \text{ V/K} = 0.347 \text{ mV/K.}$$

Problem 3

a) The results can be found in Figure 1 for the various regions:

- 1: R-MSA(s) + (R-MSA + S-IPSA)(l)
- 2: (R-MSA + S-IPSA)(l)
- 3: (R-MSA/S-IPSA)(s) + (R-MSA + S-IPSA)(l)
- 4: (R-MSA/S-IPSA)(s) + (R-MSA + S-IPSA)(l)
- 5: S-IPSA(s) + (R-MSA + S-IPSA)(l)
- 6: S-IPSA(s) + (R-IPSA + S-IPSA)(l)
- 7: (R-IPSA + S-IPSA)(l)
- 8: (R-IPSA/S-IPSA)(s) + (R-IPSA + S-IPSA)(l)
- 9: (R-IPSA/S-IPSA)(s) + (R-IPSA + S-IPSA)(l)
- 10: R-IPSA(s) + (R-IPSA + S-IPSA)(l)
- 11: (R-IPSA + R-MSA)(l)
- 12: R-MSA(s) + (R-MSA/S-IPSA)(s)
- 13: (R-MSA/S-IPSA)(s) + S-IPSA(s)
- 14: S-IPSA(s) + (S-IPSA/R-IPSA)(s)
- 15: (S-IPSA/R-IPSA)(s) + R-IPSA(s)
- 16: R-IPSA(s) + (R-IPSA + R-MSA)(l)
- 17: R-MSA(s) + (R-IPSA + R-MSA)(l)
- 18: R-IPSA(s) + R-MSA(s)

NOTE: Advised but not necessary is putting the further data in the graph (see Figure 2).

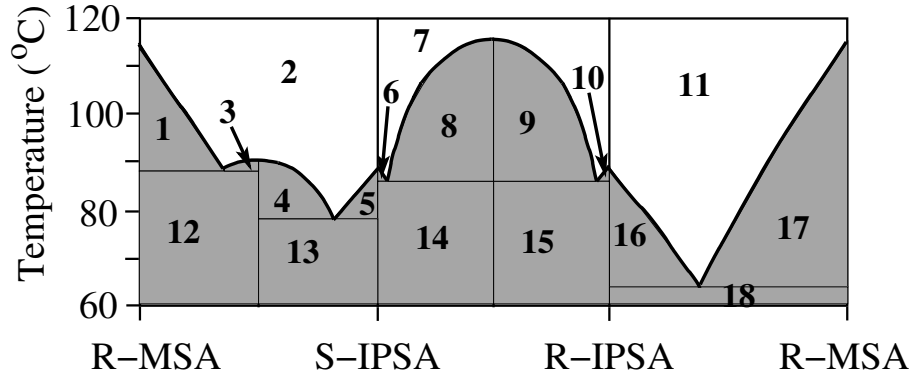


Figure 1: NOTE: All shaded areas are two-phase regions.

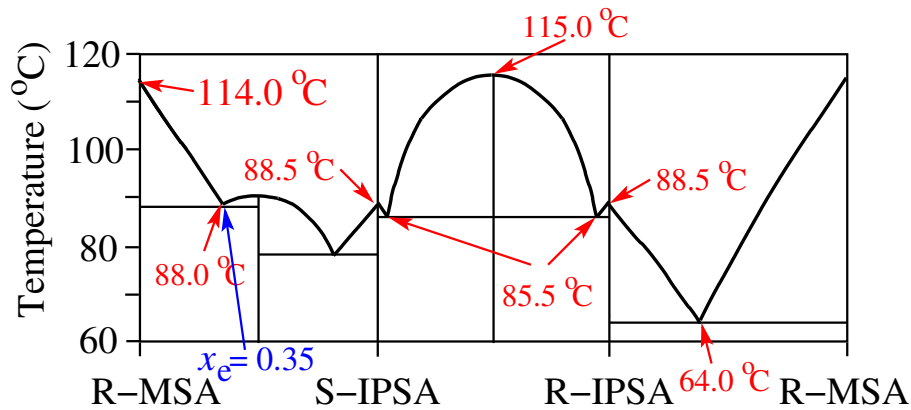


Figure 2: The further data have been added to the phase diagram.

- b) Because the phase equilibrium can be considered as linear we can use the approximate expression for freezing point depression

$$\Delta T = \left(\frac{RT^{*2}}{\Delta_{trs}H} \right) x_B = \left(\frac{R(T_{R-MSA}^*)^2}{\Delta_{fus,R-MSA}H^*} \right) x_{S-IPSA}.$$

Note: the mole fraction values run between 0 and 1 for each of the three parts of the phase diagram (R-MSA to S-IPSA, S-IPSA to R-IPSA and R-IPSA to R-MSA).

The slope of the linear line is given by

$$\frac{d\Delta T}{dx_B} = \frac{RT^{*2}}{\Delta_{fus}H^*} = \frac{114.0 - 88.0}{0.35 - 0} = 74.29 \text{ degrees K},$$

so

$$\Delta_{fus}H^* = \frac{RT^{*2}}{\frac{d\Delta T}{dx_B}} = \frac{8.314 \cdot (273.15 + 114.0)^2}{74.29} = \frac{1.246 \cdot 10^6}{74.29} = 16.77 \text{ kJ/mol}.$$

Alternatively you can use (see also problem 3d)

$$\ln x_B = \frac{\Delta_{fus}H^*}{R} \left(\frac{1}{T_{fus}^*} - \frac{1}{T} \right),$$

leading to the same result.

c)

$$\frac{n_l}{n_s} = \frac{l_s}{l_l} = \frac{1 - 0.7}{0.7 - 0.29} = \frac{0.3}{0.41} = 0.73.$$

d) Use the notation RS = R-ISPA/S-IPSA, R = R-ISPA and S = S-IPSA.

$$\mu_{\text{RS}(s)}^* = \mu_{\text{RS}(l)}^* + RT \ln a_{\text{RS}(l)}, \quad \text{so} \quad -\Delta_{\text{fus}} G_m^*(RS) = RT \ln a_{\text{RS}(l)}.$$

Because for the pure component at T_{fus}^*

$$\Delta_{\text{fus}} G_m^*(RS)(T_{\text{fus}}^*) = 0 \quad \text{it follows that} \quad \Delta_{\text{fus}} H_m^*(RS)(T_{\text{fus}}^*) - T_{\text{fus}}^* \Delta_{\text{fus}} S_m^*(RS)(T_{\text{fus}}^*) = 0,$$

so

$$\Delta_{\text{fus}} S_m^*(RS)(T_{\text{fus}}^*) = \frac{\Delta_{\text{fus}} H_m^*(RS)(T_{\text{fus}}^*)}{T_{\text{fus}}^*},$$

we find

$$-\frac{\Delta_{\text{fus}} H_m^*(RS)(T_{\text{fus}}^*)}{R} \left[\frac{1}{T} - \frac{1}{T_{\text{fus}}^*} \right] = \ln a_{\text{RS}(l)},$$

so for $T = T_e$ we find

$$-\frac{\Delta_{\text{fus}} H_m^*(RS)}{R} \left[\frac{1}{T_e} - \frac{1}{T_{\text{fus}}^*} \right] = \ln a_{\text{RS}(l)} = \ln x_{\text{RS}(l)} + \ln \gamma_{\text{RS}(l)}.$$

plugging in the values we find

$$-\frac{25.05 \cdot 10^3}{8.314} \left[\frac{1}{273.15 + 85.5} - \frac{1}{273.15 + 115} \right] = \ln(1 - 2 \cdot 0.05) + \ln \gamma_{\text{RS}(l)}, \quad \text{so} \quad \ln \gamma_{\text{RS}(l)} = -0.5331, \quad \text{and}$$

$$\gamma_{\text{RS}(l)} = 0.59.$$

Problem 4

a) For that we need the Kelvin equation

$$P = P^* \exp \left[\frac{V_m \Delta P}{RT} \right], \quad \text{with} \quad \Delta P = \frac{2\gamma}{r}$$

$$\begin{aligned} \log P^*(298 \text{ K}) &= 45.465 - \frac{4044}{298} - 13.21 \log 298 + 1.098 \cdot 10^{-7} \cdot 298 + 6.415 \cdot 10^{-6} \cdot 298^2 \\ &= 45.465 - 13.57 - 32.68 + 3.27 \cdot 10^{-5} + 0.570 = -0.215, \quad \text{so} \end{aligned}$$

$$P^*(298 \text{ K}) = 0.609 \text{ Torr} = \frac{0.609}{750} = 0.812 \text{ mbar}.$$

$$\Delta P = \frac{2\gamma}{r} = \frac{2 \cdot 43.8 \cdot 10^{-3}}{0.5 \cdot 10^{-6}} = 175.2 \cdot 10^3 \text{ Pa}.$$

$$V_m = \frac{M}{\rho} = \frac{78.13}{1.1004} \text{ cm}^3/\text{mol}.$$

Putting these values into the Kelvin equation results in

$$P = 0.812 \exp \left[\frac{71.0 \cdot 10^{-6} \cdot 175.2 \cdot 10^3}{8.314 \cdot 298} \right] = 0.816 \text{ mbar} = 0.612 \text{ Torr}.$$

b) Because $\gamma = \gamma(T_{vap}(r))$, we first determine its value

$$\gamma(T_{vap}) = G \left(1 - \frac{T_{vap}}{T_c}\right)^n = 93.26 \cdot 10^{-3} \left(1 - \frac{462.15}{726}\right)^{1.43} = 21.93 \text{ mN/m.}$$

Plugging in all values into the equation for $T'_{vap}(r)$ gives ($P = P^\ominus$)

$$T'_{vap}(0.5 \cdot 10^{-6}) = \frac{462.15}{1 - \frac{8.314 \cdot 462.15}{42.70 \cdot 10^3} \ln \left[1 + \frac{2 \cdot 21.93 \cdot 10^{-3}}{10^5 \cdot 0.5 \cdot 10^{-6}}\right]} = \frac{462.15}{1 - 0.0900 \ln 1.877} = 490 \text{ K.}$$

c) For higher temperatures T the surface tension $\gamma(T)$ will be smaller, so T'_{vap} will become closer to T_{vap} .

Furthermore $\Delta_{vap}H_m$ is in reality dependent on T ; For increasing temperatures T the vaporization enthalpy will decrease, so T'_{vap} will become closer to T_{vap} .

In practice nucleation of the gas bubble will be heterogeneous, that is, will nucleate on foreign surfaces, like dust or the wall of the vessel, resulting in a smaller effective surface tension.

d) Along the phase boundary line (the liquid-vapour coexistence line)

$$d\mu_g = d\mu_l, \quad \text{so} \quad (V_m dP - S_m dT)_g = (V_m dP - S_m dT)_l, \quad \text{or}$$

$$(S_m^g - S_m^l) dT = (V_m^g - V_m^l) dP \approx V_m^g dP.$$

Using the equation of state for a perfect gas ($PV_m = RT$) we find

$$\left(\frac{RT}{P} dP\right)_g = (S_m^g - S_m^l) dT, \quad \text{so} \quad R \left(\frac{dP}{P}\right)_g = (S_m^g - S_m^l) \frac{dT}{T}.$$

On the phase boundary line it also holds that

$$\mu_g = \mu_l, \quad \text{so} \quad H_m^g - TS_m^g = H_m^l - TS_m^l, \quad \text{or} \quad S_m^g - S_m^l = \frac{H_m^g - H_m^l}{T},$$

which we substitute in the equation for $R \left(\frac{dP}{P}\right)_g$:

$$R \left(\frac{dP}{P}\right)_g = (H_m^g - H_m^l) \frac{dT}{T^2}.$$

Integrating both sides we get (the gas in the bubble will have a higher pressure $P^\ominus + \Delta P = P^\ominus + \frac{2\gamma}{r}$; assume that γ is independent of T)

$$\int_{P^\ominus}^{P^\ominus + \frac{2\gamma}{r}} \frac{dP}{P} = \frac{(H_m^g - H_m^l)}{R} \int_{T_{vap}}^{T'_{vap}} \frac{dT}{T^2}, \quad \text{so} \quad \ln \frac{P^\ominus + \frac{2\gamma}{r}}{P^\ominus} = \frac{\Delta_{vap}H_m}{R} \left[\frac{1}{T_{vap}} - \frac{1}{T'_{vap}} \right].$$

Rearranging the expression we find

$$T'_{vap} = \frac{T_{vap}}{1 - \frac{RT_{vap}}{\Delta_{vap}H_m} \ln \left[1 + \frac{2\gamma}{P^\ominus r}\right]}.$$